AFGL-TR- 88-0067 ENVIRONMENTAL RESEARCH PAPERS, NO. 999

High Resolution Infrared Spectroscopy of Carbon Dioxide and Nitrous Oxide at Elevated Temperatures

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11 March 1988



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OPTICAL PHYSICS DIVISION PROJECT 2310

AIR FORCE GEOPHYSICS LABORATORY

HANSCOM AFB, MA 01731

REPORT	DOCUMENTATIO	N PAGE			Form Approved OMB No. 0704-0188						
a. REPORT SECURITY CLASSIFICATION Unclassified		16. RESTRICTIVE	MARKINGS		<u></u>						
2a. SECURITY CLASSIFICATION AUTHORITY		3. DISTRIBUTION	Y/AVAILABILITY (OF REPORT							
26. DECLASSIFICATION / DOWNGRADING SCHEDU	JLE	Appro Dist	ved for pub cribution un	lic rele limited.	ase;						
4 PERFORMING ORGANIZATION REPORT NUMB AFGL-TR-88-0067 ERP, No. 999	ER(S)	5. MONITORING	ORGANIZATION	REPORT NU	MBER(S)						
Air Force Geophysics Laboratory	6b OFFICE SYMBOL (If applicable) OPI		onitoring organice Geophysic		ratory						
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8c. ADDRESS (City, State, and ZIP Code)			FUNDING NUMBE								
		PROGRAM ELEMENT NO. 61102F	PROJECT NO: 2310	TASK NO	WORK UNIT ACCESSION NO 1 12						
11 TITLE (Include Security Classification)	_ 	011021	4070	<u> </u>							
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16. Supplementary notation *Stewar This research was partially su Radiance Laboratory, Utah Sta	t Radiance Labo pported under C	ontract F19	628~83-C-0	0056 wit	th Stewart 01730						
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22. NAME OF RESPONSIBLE INDIVIDUAL George A. Vanasse	W. GONCUSERS	<u> </u>	(Include Area Cod	e) 22c OF	FICE SYMBOL FGL/OPI						

Unclassified

Preface

The authors would like to thank Denise Ivaldi for her help in preparing the manuscript and Charles P. Dolan, Jr. for preparing the figures.

This work was supported by the Air Force Office of Scientific Research (AFOSR) as part of AFGL Task 2310G1.



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High Resolution Infrared Spectroscopy of Carbon Dioxide and Nitrous Oxide at Elevated Temperatures

1. INTRODUCTION

CAIR N DIOX DE

Although both 100_2 and 100_2 are minor constituents of the terrestrial atmosphere, they play leading roles in several current atmospheric problems. Both 100_2 and 100_2 are greenhouse gases, and while it is known that the atmospheric concentrations of both gases is increasing. The impact of these trends on global temperature is not yet adequately understood. In addition to being a greenhouse gas, 100_2 plays an indirect part in ozone chemistry.

In addition to the direct knowledge gained by studying the high temperature spectra of gases, these studies help increase the general understanding of the

⁽Received for publication 2 March 1988)

Hudson, R.D. and Reed, E.I. (1979) <u>The Stratosphere</u>: <u>Present and Future</u>, NASA Report 1049

^{2.} Weiss, R.F. (1981) The temporal and spatial distribution of tropospheric nitrous oxide, J. Geophys. Res., 86(C8):7185-7195.

^{3.} Wang, W.C., Yung, Y.L., Lacis, A.A., Mo, T., and Hansen, J.E. (1976) Greenhouse effects due to man-made perturbations of trace gases, Science, 194(No. 4266):685-690.

^{4.} Crutzen, P.J. (1971) Ozone production rates in an oxygen, hydrogen, nitrogen-oxide atmosphere, J. Geophys. Res., 76:7311-7327.

^{5.} World Meteorological Organization, Global Ozone Research and Monitoring Project, WMO Report No. 16 (1985).

provides of molecules. Heating a molecule makes possible the observations of molecular transitions originating from highly excited rotation-vibration states. These observations can then be used to characterize the shape of the molecular potential function at increasing distances from the minimum of the potential function.

The infrared spectra of linear molecules like CO₂ and N₂O are composed of vibration bands made up of a number of nearly equally spaced rotation lines. At high temperatures, there is a great deal of overlapping of these bands. This overlapping causes two main problems. First, the line density in the experimental spectrum becomes very high, requiring high spectral resolution. Second, this overlapping masks the regular structure of each band, complicating line assignment. To assign lines, it is very helpful to have wide continuous spectral coverage so that entire band systems can be observed. A Fourier Spectrometer is ideally suited to the study of high temperature gases since it meets both of these needs by providing high resolution over a wide spectral region.

The AFGI, high resolution Fourier Spectrometer has been used to carry out an ongoing study of atmospheric gases at elevated temperatures. This report summarizes the results that have been obtained during the present reporting period. Spectra were recorded of $\rm N_2O$ and three different isotopic species of $\rm CO_2$ in several different wavelength regions. The observed lines in the experimental spectra were then identified with individual molecular transitions for as many molecular transitions as possible (over 12,000 transitions). The final step was to use a least-squares fit to calculate new effective molecular constants.

2. THEORETICAL BACKGROUND

From a fundamental point of view, ${\rm CO}_2$ and ${\rm N}_2{\rm O}$ possess similar physical properties. Both ${\rm CO}_2$ and ${\rm N}_2{\rm O}$ are linear, triatomic molecules. They are isoelectronic and have a nuclei of nearly equal charge and mass. Only the electronic ground state of either molecule is significantly populated in the present experiment, yet its configuration plays an important role in determining the rotation-vibration energy levels. This occurs through the coupling of the electronic spin or orbital angular momentum to the motion of the nuclei. However, in the ground electronic state of both ${\rm CO}_2$ and ${\rm N}_2{\rm O}$ the net orbital and spin angular momenta of the electrons are zero, so there is no net electronic angular momentum to couple with the motion of the nuclei. There is a Fermi resonance between the bending mode, $2\nu_2$, and

^{6.} Bowens-Jenkins, P.E., Cooper, D.L., and Richards, W.G. (1985) Ab initio computation of molecular similarity, J. Phys. Chem., 89(No. 11).

the symmetric stretching mode, ν_1 , in both molecules. The Coriolis interactions and I-type doubling also occur for both molecules. The Coriolis interaction cocurs because as the molecule rotates, the asymmetric stretching mode ν_3 becomes coupled with the bending mode ν_2 . For vibrational bands where the I-type doubling interaction occurs, each rotation-vibration energy level splits into two levels; an "e" level with a symmetric wave function, and an "f" level with an antisymmetric wave function. The degeneracy of the e and the f energy levels is removed by the rotation of the molecule.

In several respects, however, ${\rm CO_2}$ differs from ${\rm N_2O}$. The main difference between the two molecules is the symmetry. ${\rm CO_2}$ has a symmetric structure O-C-O where the ${\rm N_2O}$ structure is asymmetric N-N-O. These symmetry differences greatly affect the nature of the spectra. Due to the symmetry of the ${\rm CO_2}$ molecule, transitions involving the symmetry stretch mode, ν_1 , are not dipole allowed. In addition, alternating lines are missing from rotation-vibration transitions (they have zero statistical weight). If both of the oxygen atoms in the ${\rm CO_2}$ molecule are not of equal mass, the symmetry of the molecule is broken and the character of the spectrum becomes more like that of ${\rm N_2O}$. This study covered both the symmetric isotopic species ${}^{12}{\rm C^{16}O_2}$, ${}^{13}{\rm C^{16}O_2}$, and ${}^{12}{\rm C^{18}O_2}$ and the asymmetric species ${}^{12}{\rm C^{16}O^{18}O}$, ${}^{13}{\rm C^{16}O^{18}O}$, and ${}^{13}{\rm C^{16}O^{17}O}$. The rotation-vibration energy term values, ${\rm T(v, J)}$, of a linear molecule can

The rotation-vibration energy term values, T(v, J), of a linear molecule can be expressed as a power series in J(J+1), that is

$$T(v, J) = G_v + B_v J(J+1) - D_v [J(J+1)]^2 + H_v [J(J+1)]^3 + L_v [J(J+1)]^4$$
 (1)

where G_V , B_V , and so on, are effective molecular constants. Each line in the experimental spectra corresponds to the transition between a pair of rotation-vibration states. For those bands where l-type doubling occurs (l>0) two sets of effective molecular constants are used, one for the elevels and the other for the flevels.

The notation of the vibrational states that were used for N_2^O and CO_2 is different. For N_2^O the notation was $v_1^{v_2^{l}v_3}$. For CO_2 the AFGL notation 9 was used.

^{7.} Tidwell, E.D., Plyler, E.K., and Benedict, W.S. (1960) Vibration-rotation bands of N₂O, J. Opt. Soc. Am., 50(No. 12):1243.

^{8.} Herzberg, G. (1945) Molecular Spectra and Molecular Structure, Vol. II, Van Nostran Reinhold, New York.

McClatchey, R.A., Benedict, W.S., Clough, S.A., Burch, D.E., Calfee, R.F., Fox, K., Rothman, L.S., and Garing, J.S. (1973) <u>AFCRL-Atmospheric</u> <u>Absorption Line Parameters Compilation</u>, AFCRL-TR-73-0096, AD 762904.

In the AFGL notation the vibrational states are identified by $v_1v_2Iv_3r$, where "r" is the ranking index assigned to each member of a Fermi resonating group of levels. When a state is not involved in Fermi resonance, r=1 and the AFGL notation is essentially the same as the notation used for N_2O . When Fermi resonance is present the ranking index, r, is appended to the quantum numbers of the interacting state with the highest v_i . For example, the AFGL notation for the two states 10^00 and 02^00 , which are highly mixed by Fermi resonance, is 10001 and 10002.

3. MEASUREMENTS PERFORMED

The Air Force Geophysics Laboratory high resolution interferometer was used in conjunction with a high temperature absorption cell to make the spectral measurements. The $\rm N_2O$ and $\rm CO_2$ samples that were used in the study were heated to temperatures up to 800 K. The $\rm N_2O$ spectra were taken in the 8 μm region using a $\rm N_2O$ sample of natural isotopic abundance. The $\rm CO_2$ spectra were taken in the 2.8 μm and 4.3 μm regions. Three different isotopic samples of $\rm CO_2$ were used, a sample of natural isotopic abundance, a sample enriched in $\rm ^{13}C$, and one enriched in $\rm ^{18}O$. As parts of this work have been completed, the results have been incorporated into two previous AFGL Technical Reports. $\rm ^{10,11}$ Table 1 gives a summary of where these results, including line positions, can be found for each molecule and spectral region. These results have also been published in The Journal of Molecular Spectroscopy. $\rm ^{12,13}$ In addition, the information on line position of $\rm CO_2$ have been incorporated into the 1986 edition of the AFGL HITRAN publicular database. $\rm ^{14}$

^{10.} Esplin, M.P., Sakai, H., Rothman, L.S., Vanasse, G.A., Barowy, W.M., and Huppi, R.J. (1986) <u>Carbon Dioxide Line Positions in the 2.8 and 4.3 Micron Regions at 800 Kelvin</u>, AFGL-TR-86-0046, ADA 173808.

^{11.} Barowy, W.M., Esplin, M.P., Vanasse, G.A., and Huppi, R.J. (1987)

Medium- and Long-Wave Infrared Absorption Spectra of Carbon Dioxide
and Nitrous Oxide at 800K, AFGL-TR-87-0016, ADA 179430.

^{12.} Esplin, M.P., and Rothman, L.S. (1983) Spectral measurements of high temperature isotopic carbon dioxide in the 4.3 μ m-region, J. Mol. Spectrosc., 116:351.

Esplin, M. P., and Rothman, L.S. (1986) Spectral measurements of high temperature isotopic carbon dioxide in the 4.5- and 2.8-μm regions, J. Mol. Spectrosc., 100:193.

^{14.} Rothman, L.S., Gamache, R.R., Goldman, A., Brown, L.R. Toth, R.A., Pickett, H.M., Poynter, R.L., Flaud, J.-M, Camy-Peyret, C., Barbe, A., Husson, N., Rinsland, C.P., and Smith, M.A.H. (1987) the HITRAN database: 1986 edition, Appl. Opt., 26:4058.

Table 1. Summary of Results of High Temperature Studies

Molecule	Number of Bands	Wavelength Region (μm)	Where Results Reported
$^{14}\mathrm{N_2}^{16}\mathrm{O}$	18	8	This report
$^{12}C^{16}O_{2}$	19	4.3	AFGL-TR-86-0046*
$^{13}C^{16}O_{2}$	15	4.3	AFGL-TR-86-0046*
$^{13}C^{16}O^{18}O$	5	4.4	AFGL-TR-86-0046*
$^{12}C^{18}O_{2}$	5	4.3	AFGL-TR-86-0046
$^{12}\text{C}^{16}\text{O}^{18}\text{O}$	5	4.3	AFGL-TR-86-0046*
12 _C 16 _O 18 _O	5	4.4	AFGL-TR-87-0016**
13 _C 16 _O 17 _O	1	4.4	AFGL-TR-86-0046*
$^{12}C^{16}O_{2}$	11	2.7	AFGL-TR-86-0046*
12 _C 16 _O 18 _O	10	2.8	AFGL-TR-86-0046*
¹² C ¹⁶ O ₂	2	2.8	AFGL-TR-86-0046*

^{*}ADA 173808

4. THE EXPERIMENTAL SETUP

The experimental apparatus consists of an infrared source, a high temperature absorption cell and the AFGL high resolution interferometer. A Nernst glower was used as the source of the infrared energy. The high temperature cell that was used has been described previously. The central one-meter section of this cell can be heated to 800K and is triple passed using the Pfund configuration. The total absorption path of the cell is 3.5 meters, 3 meters of uniform high

ADA 179430

^{15.} Dalton, W.S., and Sakai, H. (1980) Absorption cell for the infrared spectroscopy of heated gas, Appl. Opt., 19:2413.

temperate e and 1/4 meter on each end of the cell where the temperature drops to near ambient. Significant features of the AFGL High Resolution Interferometer include the use of cat's eye retro-reflectors, step and integrate instead of continuous carriage motion, and a digital demodulation and integration scheme.

The primary advantage of cat's eye retro-reflectors over flat mirrors is that cat's eyes are insensitive to tilt making it much easier to maintain alignment as the interferometer is scanned. A cat's eye retro-reflector also laterally displaces input and output beams making it possible to access both output beams. The optical signals from these two beams are complementary and so it is possible to use two detectors and operate them in a push-pull mode thus canceling out common mode errors. Using dual detectors also helps to reduce the effects of nonlinear detectors. The two Cu:Ge detectors of the AFGL high resolution interferometer are mounted in the same liquid helium dewar. Using only one dewar reduces cooling costs and increases convenience. It also helps match the conditions experienced by the two detectors making the common mode rejection work better.

In our apparatus the infrared beam is chopped before entering the high temperature absorption cell. The infrared signals are then detected, demodulated, and integrated digitally. The digital data system allows for fast settling time after a step, but long integration time during data taking. It is also used to compensate for small amounts of chopper litter and slight phase variations between the two complementary infrared channels.

Several components of the experimental apparatus have been reworked during this reporting period. $^{10,\,11}$ Previously, the maximum usable wavelength of the interferometer was about 7 μ m, but by installing a KBr beamsplitter and Ge:Cu detectors the usable wavelength coverage has been extended to approximately 20 μ m. The infrared source chamber was also totally rebuilt. In addition to these modifications, others are underway to increase the accuracy, reliability, and ease of use of the interferometer. These additions include a new stabilized reference laser, an improved KBr beamsplitter, a remotely operable filter wheel, and a new data system.

^{16.} Guelachvili, G. (1986) Distortion free interferograms in Fourier transform spectroscopy with nonlinear detectors, Appl. Opt., 25:4644.

4.1 Reference Laser

A new stabilized reference laser, a Laboratory for Science Model 220, has been installed in the interferometer. It has a long-term frequency drift of less than 50 kHz/day. Long-term stability is particularly important for use with a step and integrate interferometer like the AFGL high resolution interferometer where each spectral scan can require up to 15 hours. If the potential stability of the laser is to be realized, the laser must be maintained in a controlled environment. In addition to controlling the physical environment for the laser, retroreflections of the laser beam must also be controlled.

The laser head of the Laboratory for Science Model 220 laser is physically separate from the power supply and most of the other electronics. It is only the laser head that must be placed in a controlled environment. The interferometer is operated in a vacuum, but the laser must be kept at atmospheric pressure so the laser head was placed in a pressurized enclosure and cables routed to the exterior of the vacuum chamber where the rest of the electronics were located. The previous reference laser also needed to be maintained at atmospheric pressure, but due to different mechanical designs of the two lasers it was not possible to place the new laser in the old laser enclosure; thus it was necessary to design and construct a new laser enclosure. The design of this enclosure is given in Figure 1.

The long term stability of the laser is very sensitive to maintaining the laser head at a fixed temperature. To do this, we maintain the laser head packaging at an elevated temperature. This presents a problem because the pressure enclosure must provide adequate ventilation to prevent temperature buildup. The laser head can be cooled by circulating air through the enclosure if extreme caution is used to ensure that the air flow is laminar. Passing the laser beam through a turbulent airflow would introduce fluctuation in the laser beam and turbulent air around the laser head would interfere with the operation of the laser cavity length servo.

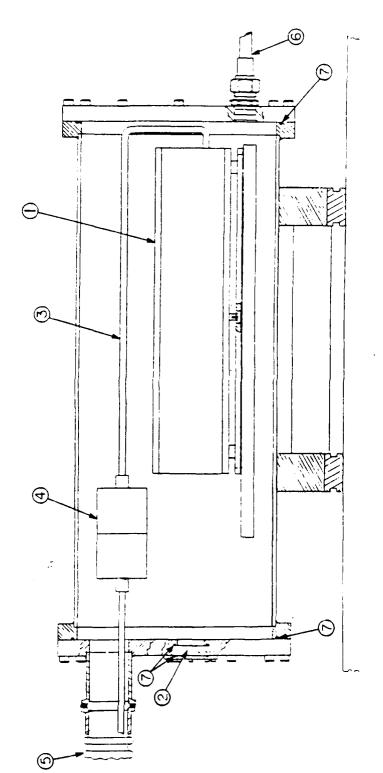


Figure 1. Reference Laser Pressure Enclosure. (1) Reference laser, (2) output window, (3) electrical cables, (4) cable connector, (5) flexible coupling to outside of vacuum chamber, (6) cooling line, and (7) ()-ring seals

An additional complication arises since the reference laser uses transverse Zeeman stabilization. This method of stabilization is very susceptible to retroreflected radiation. Careful design was needed to insure that retro-reflected energy from the interferometer did not reenter the laser. An interferometer like the AFGL High Resolution Interferometer that uses cat's eye retro-reflectors is very susceptible to backscattered radiation problems. Backscattered radiation comes primarily from the cat's eve secondary where the laser beam is brought to a focus. A small particle of dust or surface imperfection on the secondary can easily scatter considerable light back into the laser. We have found that a neutral density filter provided adequate isolation from backscattered radiation. Passing the beam through a neutral density filter reduces backscattered radiation since the light that is backscattered has to pass through the filter twice while the desired output beam only passes through the filter once. Hence, a neutral density filter with an attenuation of 10 will reduce the backscattered radiation by a factor of 100. The new laser is more powerful than the old laser, so the reduction in intensity by a factor of 10 is not a serious problem. If it proves necessary in the future, very much higher levels of isolation can be accomplished by passing the laser beam through a polarizer and a quarterwave plate.

4.2 Beamsplitter

The KBr beamsplitter used in the AFGL High Resolution Interferometer to perform the $\rm N_2O$ measurements had some deficiencies and so has now been replaced. It had been in storage for a number of years and had lost some of its flatness, although it performed satisfactorily in the longer wavelength regions used for the $\rm N_2O$ measurements. In addition, the beamsplitter coating was such that the RT product of the beamsplitter was low in the 4 μm region. The coatings on the old beamsplitter consisted of a single layer of germanium. The germanium coating was thicker on the portion of the beamsplitter used for the infrared than for the reference laser. The new beamsplitter uses different coating materials for the two regions.

The most serious problem with the old beamsplitter was that the germanium coating used for the reference laser was excessively absorbing. The absorption was high enough to make the beamsplitter appear more like a metallic than a dielectric beamsplitter. The difference between an interferometer using a metallic beamsplitter and a dielectric beamsplitter is the phase between the two beams. With a metallic beamsplitter the outputs of the two beams are in phase while with a dielectric beamsplitter they are complementary. With a dielectric beamsplitter,

^{17.} Mertz, L. (1965) Transformation in Optics, John Wiley, New York.

since the signals from the two detectors are complementary, it is possible to operate the electronics in a push-pull mode resulting in common-mode rejection. With the old KBr beamsplitter it was possible to operate the infrared detectors in a complementary manner, but not the laser reference detectors. With the new beamsplitter it will be possible to operate both the infrared and the laser reference detectors in a complementary mode.

4.3 Filter Wheel

Although a Michelson Interferometer can cover a very wide wavelength range, the photon noise produced at each spectral interval is spread through the entire spectrum. Hence, higher signal-to-noise ratios are attained if the wavelength range of the input radiation is limited with an optical filter. Since this optical filter is located inside the interferometer enclosure, it was necessary to deflate the vibration isolation pads and bring the interferometer enclosure to atmospheric pressure to change this optical filter. Since the AFGL High Resolution Interferometer is a slow scanning instrument (up to 15 hours), the long-term stability of the instrument is extremely critical. More consistent results are obtained if the interferometer is allowed to equilibrate for several days after having been opened. Hence changing the optical filter resulted in several days of lost opportunity to take data. To get around this problem a six position filter wheel has been installed in the interferometer vacuum enclosure that can be operated from outside of the vacuum tank. The design of this filter wheel is shown in Figure 2.

4.4 Data Aquisition System

A new data system is currently being implemented using an IBM AT compatible computer to replace the old system which was based on an outdated PDP 8/e computer. With the PDP 8/e system, only data acquisition and control of the interferometer were performed locally, and all subsequent processing of the data was performed using a Control Data mainframe computer. With the new system, much more of the processing of the data will be possible locally. The ability to perform "quick-look" checks on the data before transferring it to the mainframe will be particularly valuable. Given the rapid progress in the computer industry, in the near future it should be possible to obtain a microcomputer with sufficiently high performance to perform the entire data processing on a microcomputer.

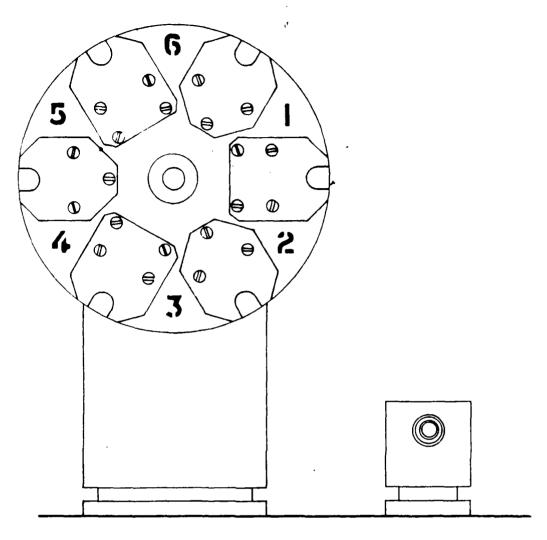


Figure 2. Design of Optical Filter Wheel

The main advantages of using an IBM PC for an interferometer controller are low costs and the flexibility made possible by the large number of available hardware and software options. However in many ways, the PDP 8/e was better suited for use as a programmable controller than an IBM PC computer. The PDP 8/e is a synchronous machine with a very simple non-intrusive operating system. With the PDP 8/e, the computer is never servicing an interrupt when a time critical operation is in progress. Also since the PDP 8/e is synchronous, the time required to perform each instruction is always the same, making it possible to perform timing by using the speed of the computer instructions themselves. With an IBM PC on the other hand, since the time required to perform a given operation is not constant, the computer cannot be counted on to perform time critical operations without the aid of additional hardware. (In principle it is possible to turn off

the interrupt on the IBM PC, but since the operating system and the hardware of the computer are very closely integrated this is not a very satisfactory solution.) These disadvantages are largely overcome by the manufacturers of the data acquisition systems supplying complete data acquisition subsystems instead of just simple analog to digital converters. For example, the MetraByte Dash-16 board used in the new data acquisition system includes, in addition to the analog to digital converters, timing and logic circuits as well as the necessary software drives to use them.

The new data acquisition and control software is written in FORTRAN 77, in contrast to the old PDP 8/e system which was written in absolute loaded assembly language. Assembly language programs in general tend to use the hardware more efficiently, but they are also much more dependent on the details of the hardware. The reason we have been using a PDP 8/e long after it has become obsolute is because of the high manpower cost required to rewrite the data acquisition and control program to make use of new hardware. It should be relatively easy to modify the new system to keep up with advances in hardware. The ease of programming in a higher level language and the availability of commercially available graphics and mathematical software, also makes it practical to write software which is considerably easier to use and much more capable.

5. N₂O MEASUREMENTS

Spectra were measured at several different temperatures and pressures to facilitate the line assignment process and to maximize the number of spectral lines measured under optimum conditions. After checking each of these spectra for consistency, they were co-added to obtain one spectrum for each temperature and pressure. The experimental conditions under which spectra were measured and the number of experimental spectra taken are listed in Table 2. Measurements of the spectrum of the empty absorption cell were interspersed between the $\rm N_2O$ spectra. These empty cell measurements were used to determine the 100 percent transmission levels. The maximum retardation of the interferometer for all spectra was 83 cm, resulting in a resolution of 0.006 cm⁻¹.

Table 2. Experimental Conditions for N_2^{O} , $\delta \sigma = 0.006 \text{ cm}^{-1}$

Number of Spectra	Temperature	Pressure
2	300K	1.0 Torr
2	473 K	2.3 Torr
2	473 K	9.0 Torr
2 -	800K	4.0 Torr
2	800K	15. Torr

Additional spectra not included in Table 2 were the first obtained at $800\,\mathrm{K}$; however, few absorption features were observed. As these were spectra of the first samples run at $800\,\mathrm{K}$, it appears plausible that the high temperature $N_2\mathrm{O}$ was reacting with the walls of the cell. Residue from the initial reactions prevented further loss of $N_2\mathrm{O}$, enabling successful recording of the interferograms that followed.

The spectra were calibrated using an internal calibration technique. Toth has published a paper 18 in which he reports the analysis of room temperature $\rm N_2O$ spectra that were taken using the high resolution Fourier spectrometer located in the McMath solar telescope facility at the Kitt Peak National Observatory. Since his data were taken at room temperature the lines that he observed do not extend to as high rotational states (J values) as in this work. The calibration was performed by adjusting the wavelength scale of the observed spectra until, on the average, the observed line positions for the low J lines matched the values obtained by Toth.

During the calibration process it was noted that there were some systematic shifts between the positions of strong and weak lines in spectra taken at high temperatures. This effect is presently being investigated, but is probably due to non-uniform illumination of the detectors causing a slightly asymmetric instrumental line shape. Evidence to support this conclusion is that the quality of the infrared beam has been observed to be much poorer at the higher temperatures than at room temperature. The primary cause of this beam degradation was probably due to the

^{18.} Toth, R.A. (1986) Frequencies of N₂O in the 1100 to 1440 cm⁻¹ region, J. Opt. Soc. Am., 3:1263.

distortion of the mirrors in the absorption cell as the cell was heated. Even for the low J lines of the ν_1 fundamental where this effect was most noticeable, the error was only 0.0007 cm⁻¹.

5.1 Treatment of the Data

The line assignments for each band were performed by starting at low J, where Toth's molecular constants were valid, and working to high J. After the line assignments had been made, data from all the different temperatures and pressures were combined into a single data set and a final weighted least-squares-fit was performed to obtain new effective molecular constants. Although over 4100 lines were identified in the experimental spectrum, only 3454 were used in the least-squares-fit, due to line merging problems. Many of the remaining lines were slightly affected by the presence of close-by spectral lines. These slightly merged lines were included in the least-squares fit, but with reduced weighting. Each band was fit independently without making any attempt to combine the information from the various bands into a single global self-consistent set of energy levels for the N₂O molecule.

In order that this weighted least-squares-fitting procedure could be used, it was necessary that an estimate of the uncertainty of each experimental line be made. The weight assigned each spectral line was the reciprocal of the expected uncertainty squared. The factors that went into calculating the expected uncertainty of each line were: the random experimental noise in the spectrum, line asymmetry, abnormal width of spectral lines, and inconsistencies of line positions compared to other lines in the same band. The total uncertainty for each line was defined as the square root of the sum of squares of the individual uncertainties. Further details of the methods used to determine the weights for the least-squares-fits were described in a previous report. 10

The most noticeable effect of using a weight for each spectral line was to substantially reduce the uncertainty in the spectroscopic constants as predicted by the least-squares-fitting program. However, the spectral line positions calculated using the resulting constants were found to be quite insensitive to the values of the weights chosen. This indicates that the effects of line merging on the position of spectral lines were essentially random for the high temperature spectra considered in this study.

5.2 Data Analysis Software

Even though the data analysis software had been used in analyzing asymmetric species of CO_2 that have similar structure to $\mathrm{N}_2\mathrm{O}$ it was still necessary to modify the software to take into account that $\mathrm{N}_2\mathrm{O}$ is less "harmonic oscillator like" than CO_2 . The energy expansion in terms of G, B, D, and H doesn't work as well, so it was necessary to include an additional term, L. This is partly due to the fact that Coriolis perturbation plays a larger role with $\mathrm{N}_2\mathrm{O}$ than it does with CO_2 . Adding L's to the fitting program made it necessary to change the format of the molecular constant data files and the formats of all the subsequent programs that use these molecular constant files.

In working with the least-squares-fitting program it became apparent that because of numerical instabilities adding L's was more involved than just adding an additional term to the fitting function. The relative size of the constants that were being fitted ranged from the order of 10^3 for G to 10^{-18} for L. This large variation in the size of the fitted parameters created problems with numerical roundoff in the inversion of the matrix that was used to obtain least-squares-fit to molecular constants. It was found that the numerical stability of the molecular constants was considerably improved by scaling the constants and including their order of magnitude into the fitting equation. Therefore in the numerical matrix inversion and in the determination of the constants all the constants were of nearly the same magnitude. If it would have been necessary, additional least-squares techniques could have been used to reduce the numerical instabilities further.

In addition to changing the software to include L's, several other changes were made to make the software easier to modify in the future. One of these changes was to convert the programs from Control Data Corporation FORTRAN IV to generic FORTRAN 77. Converting to FORTRAN 77 also eliminated the inconvenience of being tied to a single computer.

There are many advantages in using more than one computer to perform data analysis. The computer work was performed partly on an IBM-PC compatible and partly with a Control Data Corporation (CDC) mainframe computer. Due to the more convenient operating environment and better editors, it was found to be more productive to perform program editing as well as much of the program development on an IBM-PC compatible. However, most of the actual data analysis was performed using the CDC mainframe due to faster execution speeds and larger disk storage size. An additional advantage of using two computers, with different word sizes and which handle floating point calculation differently, is the ability to quickly detect numerical roundoff in the algorithms that are being used. The same calculation can be run on both computers and the results compared, thus checking for numerical roundoff.

5.3 Results and Discussion

The 18 rotation-vibration bands of $N_{\phi}O$ for which molecular constants were obtained are indicated on the energy level diagram of Figure 3. The range of P and R lines used in the least-squares fits, the total number of lines, and the rms error for each band is given in Table 3. The effective molecular constants which were obtained are given in Table 4. The line position, observed minus calculated, and expected uncertainty of each line used in the least-squares-fit are given in the appendix. These constants are effective molecular constants and so should not be expected to accurately represent the internal structure of the $\mathrm{N}_{2}\mathrm{O}$ molecule. The purpose of these effective constants is to provide a means of reproducing, within the experimental accuracy, the position of spectral lines over the range of J values covered by the measurements (see Talle 3). There are a great many interactions between different vibrational states for the $\mathrm{N}_2\mathrm{O}$ molecule. The effects of these interactions are accounted for by allowing the different effective molecular constants to float freely in the least-squares-fit of each band. Molecular constants obtained in this manner are not self-consistent. For example, the molecular constants obtained for the vibrational state 1110 from the 1110 \leftarrow 0110 band are not consistent with those obtained from the 1310 \leftarrow 1110 band.

An interesting observation that becomes apparent from studying both CO_2 and $\mathrm{N}_2\mathrm{O}$ is that the interactions between levels not involved in Fermi resonances are much stronger for $\mathrm{N}_2\mathrm{O}$ than for CO_2 . These interactions result in $\mathrm{N}_2\mathrm{O}$ being less "harmonic oscillator like" than CO_2 . The effects of these interactions can be seen in Table 4 in the magnitude of the inconsistencies of effective molecular constants determined from different bands and from the extremely large values for L's which are obtained for some bands.

Each band was fit twice, once using L's and once without. The spectroscopic constants (L' for the upper state and L'' for the lower state) were included in the final least-squares fit only when their inclusion markedly improved the quality of the fit (a reduction in the rms error of more than 20 percent) and the uncertainties in L were smaller than the value of L for both the upper and the lower states. Occasionally, an exception was made for bands where l-type doubling was present (bands where l > 0). If the elevels indicated the need of an L and the flevels did not, for consistency L's were used for both sets of levels.

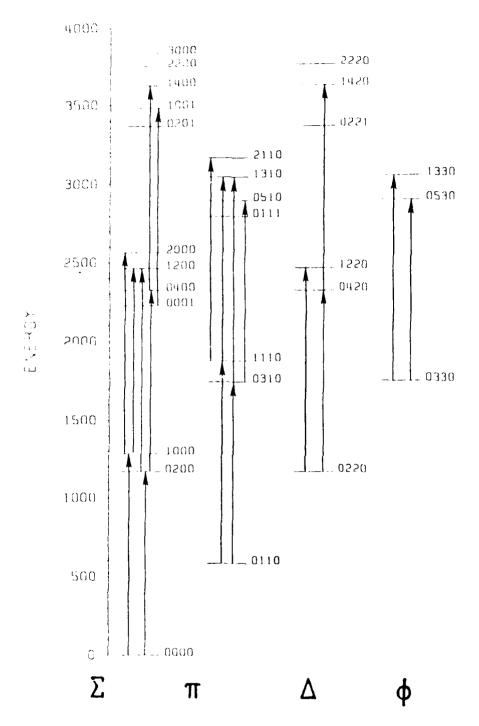


Figure 3. N_2O Energy Level Diagram Showing Observed Rotation-Vibration Bands

Table 3. Observed $\mathrm{N}_2\mathrm{O}$ Bands

		Band Center	Range of	Number	RMS Error
Transi	tion	(cm ⁻¹)	Measurement	of Lines	x10° cm ⁻¹
1000	0000	1284.9027	P(105)-R(104)	205	3
1110e	0110e	1291.4978	P(90)-R(100)	181	3
1110f	0110f	1.291.4978	P(94)-R(95)	176	3
0200	0000	1168.1319	P(83)-R(82)	156	3
0310e	0110e	1160.2973	P(73)-R(73)	137	4
0310f	0110f	1160.2973	P(75)-R(66)	132	3
1200	0200	1293.8641	P(89)-R(88)	162	3
1200	1000	1177.0927	P(64)-R(62)	104	5
1220e	0220e	1297.0542	P(83)-R(83)	155	7
1220£	0220f	1297.0542	P(89)-R(86)	147	4
2000	1000	1278.4359	P(93)-R(90)	157	3
0400	0200	1154.4403	P(63)-R(68)	119	5
0420e	0220e	1153.3767	P(59)-R(59)	101	9
0420£	0220£	1153.3767	P(65)-R(64)	102	7
1310e	0310e	1297.1481	P(83)-R(84)	125	5
1310f	0310f	1297.1481	P(86)-R(80)	124	5
1310e	1110e	1165.9488	P(58)-R(54)	60	9
1310f	1110f	1165.9488	P(60)-R(55)	61	8
1330e	0330e	1301.8082	P(86)-R(81)	136	8
1330f	0330£	1301.8082	P(81)-R(81)	142	9
053 0e	0330e	1147.1321	P(53)-R(59)	74	13
0530£	0330f	1147.1321	P(50)-R(60)	74	14
2110e	1110e	1285.5881	P(86)-R(84)	130	5
2110f	1110f	1285.5881	P(84)-R(77)	124	4
1001	0001	1257.0628	P(77)-R(68)	100	6
1400	0400	1298.3692	P(57)-R(64)	85	7
1420e	0420e	1300.4682	P(59)-R(65)	87	9
1420£	0420£	1300.4682	P(75)-R(64)	98	9

Table 4. Effective Molecular Constants (cm⁻¹)

tion 6' - 6" B' D'X10' H'X10 ¹ L'X10 ¹ B" D"X10' H'X10 ¹ L'X10 ¹ B" D'X10' L'X10' L'X		1			,		0 -			-	- ا
100 1284.9027 141725521 1.72394 1.655 1.499 1.75890 1.75990	Trans	iti		В′	D'x10'	H'X10 ¹³	L'x10 ¹⁸	В"	D"x10'	H"X1013	L"x1018
000 1284,9027 -1725521 1.72394 .6975 7.65 .41901092 1.75890 5596 2.2.7.10e 1106 01106 1291.4978 .41746590 1.75252 1.4499 .41971895 1.78690 0300 210 01106 1168.1319 .41937171 1.71522 2.41499 .419971895 1.79178 1457 200 01106 1160.2973 .42198026 2.49036 28.5104 10.11 .75136 -1.4577 51.66 200 0100 1160.2973 .42108035 2.17418 -4.1519 .41997131 1.76852 -1.7303 200 1000 127.0824 .41165296 2.47089 2.41519 .41997131 1.76872 -2.0909 200 1000 127.0824 .4116508 2.41859 2.41869 2.4114 -4.1519 .41997131 1.76872 -2.0909 200 1000 127.0824 .4185021 1.74873 1.74873 1.7003 -2.1403 -2.1413											
100 01100 1291.4978 44746590 1.75252 1.4499 1.4917895 1.78669 1.0300 1.201.4978 441837713 1.715252 2.4416 1.41991241 1.79178 -1457 1.79178 1.41991246 1.291.4978 1.4199266 2.49036 2.49036 2.49136 2.4116 1.41991241 1.75136 -1.457 1.75136 -1.457 1.6000 1.600.2973 441958062 2.10895 8.4733 441997131 1.75136 -1.7513	00	00	284.902	172552	239	697	9.	190109	.7589	559	•
110f 0110f 1291.4978 .41837171 1.71922 2.4116 .4199648 1.75178 -11457 51.65 200 0000 1168.1319 .41992266 2.49936 2.85104 10.11 .14991760 1.76186 -4.5817 51.65 310c 0110c 1160.2973 .41958062 2.08957 8.4732 9.0163 -4.5889 -2.0909 310c 0110c 1160.2973 .4108035 2.17418 -4.1519 .41997171 1.76882 -2.0909 200 1297.0642 .41104698 2.47089 21.4166 -42012475 1.9706 -2.0909 200 127.0826 1297.0642 .41652958 1.2664 -2.3114 -42012475 1.9706 -2.943 -2.943 200 1200 1200 1200 1.74403 .4260211 1.6094 -4.1755107 1.9102475 1.9106 -2.943 -2.943 200 1200 1200 1.74403 1.56098 -4.177717-122771 -4.991353 2.6484	1.10	110	291.497	174659	.7525	.449		4191789	.7860	30	
200 0000 1168.1319 41992266 2.49036 2.8.5104 10.11 .41991268 -4.5817 51.66 3100 01100e 1160.2973 44198062 2.08957 8.4733 .9.91171 1.75888 -2.0909 3100 01100e 1160.2973 42108035 2.17418 -4.1519 .9.4199108 1.76882 -2.0909 200 1200 1297.0542 441852958 1.20564 -26.314 .42012475 1.970617 -3.7088 200 1000 1297.0542 441852958 1.20564 -26.314 .42012475 1.19706 -29.6815 200 1000 1297.0542 441852958 1.74873 1.5077 42012475 1.18783 -2.243 200 1000 1294.0403 4206288 1.74873 1.5077 42012475 1.18783 -2.243 400 200 1154.4403 4206288 1.74873 1.5077 42012475 1.1878 -2.243 400 200 1154.4403 <td< td=""><td>7.10</td><td>110</td><td>291.497</td><td>183717</td><td>.7192</td><td>.411</td><td></td><td>199684</td><td>.7917</td><td>.145</td><td></td></td<>	7.10	110	291.497	183717	.7192	.411		199684	.7917	.145	
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420f 0220f 1153.3767 .42077256 2.15646 3.1991 .42012696 1.81663 8965 310e 0310e 1297.1481 .41775533 2.08241 8.4884 .41958100 2.10173 9.3416 310f 0310f 1297.1481 .41937526 2.15487 2.6078 .42107864 2.19213 .8100 310f 110f 1165.9488 .41775211 1.99684 -5.2920 .41745936 1.65249 -16.6968 310f 1110f 1165.9488 .41937305 2.16162 1.9434 .41837419 1.76630 8.9529 310f 1301.8082 .41911053 1.59673 -7.1452 -259.63 .42066992 1.63249 -16.6968 330f 0330e 1301.8082 .41911053 1.55673 -7.1452 -259.63 .42066992 1.63249 -16.6968 330f 13477.1321 .42123301 1.62107 -26.66693 1.62896 8.65539 1.69896 1.5053 110f 11285.5881	420	220	153.376	207725	1787	48.621		201269	.1778	37.140	
310e0310e1297.1481.417755332.082418.4884.419581002.101739.3416310f0310f1297.1481.419375262.154872.6078.421078642.19213.8100310e1110e1165.9488.417752111.99684-5.2920.417459361.65249-16.6968310e1110f1165.9488.419373052.161621.9434.418374191.766308.9529330e0330e1301.8082.419110531.596737.1463155.39.420669921.63249-9.4152-219.4330f0330f1147.1321.421233011.62107-26.0673.420676931.698968.6457530f0330f1147.1321.421233011.62107-2.8387.417463021.748021.2033110e1110f1285.5881.416918911.582896.2603.417463021.717171.91020011257.0628.413775881.64108-9.8693.415549411.67389-11.738840004001298.3692.418790333.87543136.8809.42077124.5553463.5614-2987.5420f0420f1300.4682.419034011.84591-124.65491599.29.420771241.89174-116.40741438.9	420	220	153.376	207725	.1564	.199		4201269	.8166	.896	
310f0310f1297.1481.419375262.154872.6078.421078642.19213.8100310e1165.9488.417752111.99684-5.2920.417459361.65249-16.6968310f1110f1165.9488.419373052.161621.9434.418374191.766308.9529330e0330e1301.8082.419110531.59673-7.1452-259.63.420669921.63249-9.4152-239.41330f0330f1301.8082.419110531.596737.1403155.39.420669921.63249-6.6500157.4530e0330f1147.1321.421233011.62107-26.0673.420676931.698968.6457530f0330f1147.1321.421233011.6210722.8387.420676931.698961.5053110f1110f1285.5881.415844271.680673.9219.417463021.748021.2033110f11257.0628.413775881.64108-9.8693.4206759411.67389-11.738840004001298.3692.418790333.87543136.8809.420625384.04056163.4596420f0420e1300.4682.419034011.84591-124.65491599.29.420771241.89174-116.40741438.9	310	310	297.148	177553	.0824	.488		4195810	.1017	.341	
310e1165.9488.417752111.99684-5.2920.417459361.65249-16.6968310f1110f1165.9488.419373052.161621.9434.418374191.766308.9529330e0330e1301.8082.419110531.59673-7.1452-259.63.420669921.63249-9.4152-239.41330f0330f1147.1321.421233011.62107-26.0673.420676931.698968.6457530e0330f1147.1321.421233011.62107-22.8387.420676931.698968.6457530f0330f1147.1321.415844271.680673.9219.417463021.748021.2033110f1110f1285.5881.416918911.582896.2603.418372921.717171.910200100011298.3692.418790333.87543136.8809.420625384.04056163.4596420e0420e1300.4682.41903401.7168892.6146-2982.78.420771241.89174-116.40741438.99	310	310	297.148	193752	.1548	.607		4210786	.1921	810	
310f1110f1165.9488.419373052.161621.9434.418374191.766308.9529330e0330e1301.8082.419110531.59673-7.1452-259.63.420669921.63249-9.4152-239.4152330f0330f1301.8082.419110531.596737.1403155.39.420669921.632496.6500157.4530e0330f1147.1321.421233011.62107-26.0673.420676931.698968.6457530f0330f1147.1321.415844271.680673.9219.417463021.748021.2033110f1110f1285.5881.416918911.582896.2603.418372921.717171.910200100011257.0628.413775881.64108-9.8693.420625384.04056163.4596420e04001298.3692.41903401.7168892.6146-2982.78.42077124.5553463.5614-2987.5420f0420f1300.4682.419034011.84591-124.65491599.29.420771241.89174-116.40741438.99	310	110	165.948	177521	.9968	5.292		174593	.6524	6.696	
330e0330e1301.8082.419110531.59673-7.1452-259.63.420669921.63249-9.4152-239.4330f0330f1301.8082.419110531.596737.1403155.39.420669921.632496.6500157.4530e0330f1147.1321.421233011.62107-26.0673.420676931.698968.6457530f0330f1147.1321.421233011.6210722.8387.420676931.698961.5053110e11285.5881.415844271.680673.9219.417463021.717171.9102110f11285.5881.416918911.582896.2603.418372921.717171.9102001100011257.0628.413775881.64108-9.8693.415549411.67389-11.738840004001298.3692.418790333.87543136.8809.420625384.04056163.4596420e0420e1300.4682.419034011.84591-124.65491599.29.420771241.89174-116.40741438.9	310	110	165.948	193730	.1616	.943		183741	.7663	.952	
330f0330f1301.8082.419110531.596737.1403155.39.420669921.632496.6500157.4530e0330e1147.1321.421233011.62107-26.0673.420676931.698968.6457530f0330f1147.1321.421233011.6210722.8387.420676931.698961.5053110e1110e1285.5881.415844271.680673.9219.417463021.748021.2033110f1110f1285.5881.416918911.582896.2603.418372921.717171.910200100011257.0628.413775881.64108-9.8693.42549411.67389-11.738840004001298.3692.418790333.87543136.8809.420625384.04056163.4596420e042077124.5553463.5614-2987.5420f0420771241300.46\$2.419034011.84591-124.65491599.29.420771241.89174-116.40741438.9	330	330	301.808	191105	.5967	7.1452	259.6	206699	.6324	9.4152	239.4
530e0330e1147.1321.421233011.62107-26.0673.420676931.698968.6457530f0330f1147.1321.421233011.6210722.8387.420676931.698961.5053110e1110e1285.5881.415844271.680673.9219.417463021.7748021.2033110f1110f1285.5881.416918911.582896.2603.418372921.717171.910200100011257.0628.413775881.64108-9.8693.415549411.67389-11.738840004001298.3692.418790333.87543136.8809.420625384.04056163.4596420e042077124.5553463.5614-2987.5420f0420771241.89174-116.40741438.9	330	330	301.808	191105	.5967	.140	55.3	206699	.6324	.650	57.4
530f 0330f 1147.1321 .42123301 1.62107 22.8387 .42067693 1.69896 1.5053 110e 11285.5881 .41584427 1.68067 3.9219 .41746302 1.71717 1.9102 110f 11285.5881 .41691891 1.58289 6.2603 .41837292 1.71717 1.9102 001 1257.0628 .41377588 1.64108 -9.8693 .41554941 1.67389 -11.7388 400 0400 1298.3692 .41879033 3.87543 136.8809 .42062538 4.04056 163.4596 420e 042077124 .55534 63.5614-2987.5 420f 042077124 1.89174-116.4074 1438.9	530	330	147.132	212330	.6210	26.067		206769	.6989	.645	
110e11285.5881.415844271.680673.9219.417463021.748021.2033110f11285.5881.416918911.582896.2603.418372921.717171.91020011257.0628.413775881.64108-9.8693.415549411.67389-11.738840004001298.3692.418790333.87543136.8809.420625384.04056163.4596420e0420e1300.4682.41903401.7168892.6146-2982.78.42077124.5553463.5614-2987.5420f0420f1300.4682.419034011.84591-124.65491599.29.420771241.89174-116.40741438.9	530	330	147.132	212330	.6210	2.838		206769	.6989	.505	
110f 1110f 1285.5881 .41691891 1.58289 6.2603 .41837292 1.71717 1.9102 001 1257.0628 .41377588 1.64108 -9.8693 .41554941 1.67389 -11.7388 400 0400 1298.3692 .41879033 3.87543 136.8809 .42062538 4.04056 163.4596 420e 0420e 1300.4682 .41903401 .71688 92.6146-2982.78 .42077124 .55534 63.5614-2987.5 420f 0420f 1300.4682 .41903401 1.84591-124.6549 1599.29 .42077124 1.89174-116.4074 1438.9	110	110	285.588	158442	.6806	.921		174630	.7480	.203	
001 0001 1257.0628 .41377588 1.64108 -9.8693 .41554941 1.67389 -11.7388 400 0400 1298.3692 .41879033 3.87543 136.8809 .42062538 4.04056 163.4596 420e 0420e 1300.4682 .41903401 .71688 92.6146-2982.78 .42077124 .55534 63.5614-2987.5 420f 0420f 1300.46\ddots .41903401 1.84591-124.6549 1599.29 .42077124 1.89174-116.4074 1438.9	110	110	285.588	169189	.5828	.260		183729	.7171	.910	
400 0400 1298.3692 .41879033 3.87543 136.8809 .42062538 4.04056 163.4596 420e 0420e 1300.4682 .41903401 .71688 92.6146-2982.78 .42077124 .55534 63.5614-2987.5 420f 0420f 1300.46\$2 .41903401 1.84591-124.6549 1599.29 .42077124 1.89174-116.4074 1438.9	00	00	257.062	137758	.6410	9.869		155494	.6738	11.738	
420e 0420e 1300.4682 .41903401 .71688 92.6146-2982.78 .42077124 .55534 63.5614-2987.5 420f 0420f 1300.46&2 .41903401 1.84591-124.6549 1599.29 .42077124 1.89174-116.4074 1438.9	40	40	298.369	187903	.8754	36.880		206253	.0405	63.459	
420f 0420f 1300.46\$2 .41903401 1.84591-124.6549 1599.29 .42077124 1.89174-116.4074 1438.9	420	420	300.468	190340	168	2.6146-	982.7	207712	553	3.5614-	987.5
	420	420	300.468	190340	84591	124.6549	599.2	207712	9174	116.4074	438.9

One of the most common ways to deal with small interactions between rotation-vibration states is through the use of contact transformations. Assuming that this technique is valid, the two sets of spectroscopic constants that occur for bands where I-type doubling occurs are not independent. For these bands, several of the spectroscopic constants for the e and the fisets of levels should be constrained to be equal. This interdependence between sets of rotational constants is a function of I. When I=1, the vibrational term values G_e and G_f should be constrained to be equal. When I=2, in addition to having G_e equal to G_f , the rotational constants B_e and B_f should be equal. When I=3, then $D_e=D_f$, and so on. These constraints seemed to work very well for the CO_2 bands considered in this study, but for N_2O they seemed to be causing some problems with the fits. One of the things that will be investigated as part of future efforts is whether these constraints should be dropped for N_2O .

The residual to the least-squares-fits are plotted in Figures 4-8 for several rotation-vibration bands. The $\rm N_2O$ line positions calculated using the molecular constants reported by other researchers are also indicated on these same plots. For low J lines, the $\rm N_2O$ positions of both Guelachvili and Toth are within the experimental error of the results presented in this report, but at higher J the measurements start to diverge. This result is not surprising since both Guelachvili and Toth used room temperature absorption cells. Of the two measurements Toth's is the most recent and covers a larger number of rotation-vibration bands. The line positions reported by Toth are also in better agreement with the values reported in this study than are Guelachvili's line positions. Although Joth's low J values are very accurate, they cannot be used to extrapolate the position of the high J lines observed in this study. For example if Toth's constants were used to predict the position of the high J lines, the error would sometimes be as large as 0.05 cm⁻¹ (see Figure 8).

^{19.} Amat, G., and Nielsen, H. H. (1958) Vibrational I-type doubling and I-type resonance in linear polyatomic molecules, J. Mol. Spectrosc., 2:152.

^{20.} Guelachvili, G. (1982) Absolute N₂O wavenumbers between 1118 and 1343 cm⁻¹ by Fourier transform spectroscopy, <u>Can. J. Phys.</u>, <u>60</u>:1334.

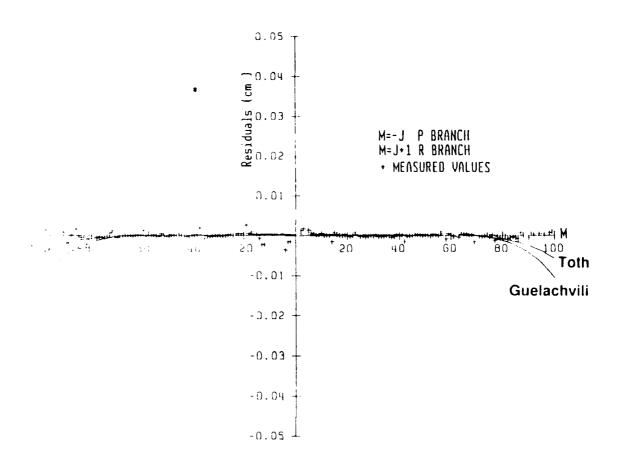


Figure 4. Comparison of Measured Line Positions With Those Computed Using Toth's and Guelachvili's Constants for the 1110e + 0110e Band

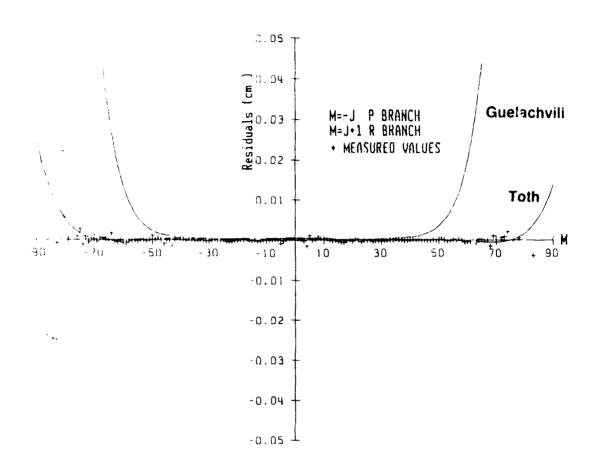


Figure 5. Comparison for the 0200 ← 0000 Band

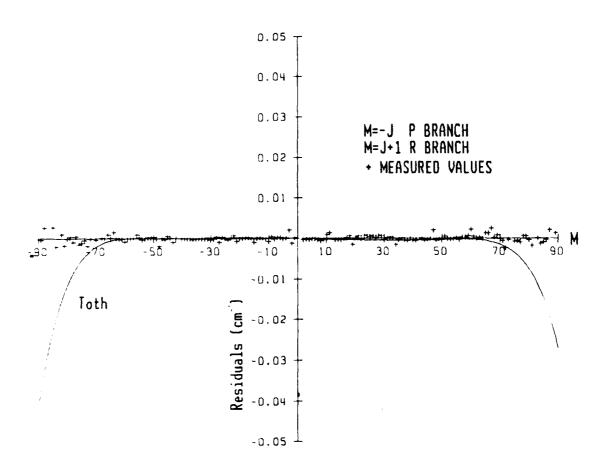


Figure 6. Comparison for the 1200 ← 0200 Band

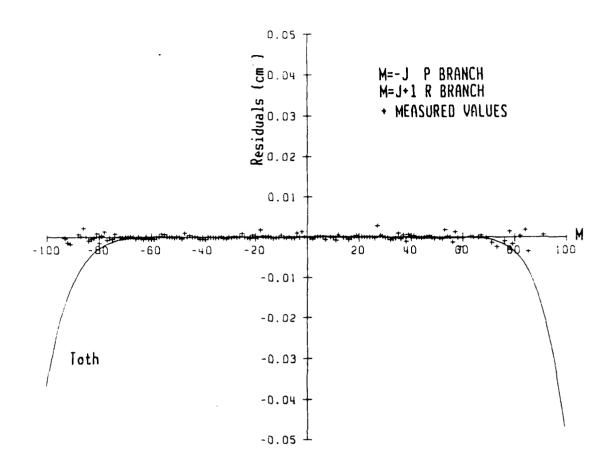


Figure 7. Comparison for the 2000 ← 1000 Band

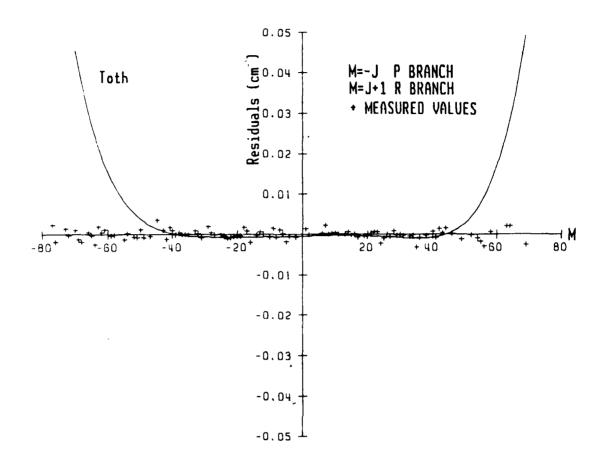


Figure 8. Comparison for the 1001 ← 0001 Band

6. CONCLUSION

During the period covered by this report an extensive study has been made on CO_2 and $\mathrm{N}_2\mathrm{O}$ using the AFGL High Resolution Interferometer in conjunction with a high temperature absorption cell. This study includes the identification of over 11,000 CO_2 lines belonging to 78 different rotation-vibration bands and over 4100 $\mathrm{N}_2\mathrm{O}$ lines belonging to 18 bands. Many of the high J lines for these bands have not been observed previously. A weighted least-squares-fit technique was then used to obtain effective molecular constants for each of these bands. These effective molecular constants predict the position of spectral lines originating from excited rotational states with an accuracy considerably greater than those previously available. The CO_2 data have already been incorporated into the 1986 edition of the AFGL HITRAN molecular database. The $\mathrm{N}_2\mathrm{O}$ data as well as the CO_2 data will be incorporated into an AFGL high temperature database that is presently being compiled.

In addition to providing data for the AFGL databases this extensive data makes it possible to come to several other interesting conclusions. Molecular constants obtained using room temperature CO_2 and $\mathrm{N}_2\mathrm{O}$ are not adequate for predicting the position of spectral lines observed at high temperatures even when the room temperature measurements are extremely accurate. The interaction between levels not involved in Fermi resonances are much stronger for $\mathrm{N}_2\mathrm{O}$ than for CO_2 . These interactions make $\mathrm{N}_2\mathrm{O}$ less "harmonic oscillator like" than CO_2 .

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Appendix A

N₂0 Line Position Listing

		10	- 00	0000) 446	1 1			2000		000 446	1 1		1 1
J.	ايما	obs	:	Unc	R Obs	0-C	Unc	sq0 d	J-0			0-c	Unc	ָר
0					285.737	! 	4			 	279.267	-1	2	0
-	28	65	2	4	286.568		4	277.601	7	٣	280.094	-4	~	1
7	28	224	17	11	287.396		٣	276.765	14	13	280.919	0	3	7
٣	28	11	-2	7	288.219		7	275.922		ო	281.741	7	3	٣
4	28	29	0	٣	289.040		e.	275.079	11	12	282.559	7	٣	4
2	1280.	9119	0	3	1289.8570	0	٣	1274.2304	- 1	7	1283.3735	0	7	വ
9	27	21	-2	٣	290.670	'	ᠬ	273.379	-2	7	284.184	-1	വ	9
7	278	63	-1	m	291.480		6	272.525	7	m				7
80	27	100	-1	٣	292.286	'	ব	271.667	7	7	285.797	,	2	80
6	27	234	-3	3	293.089		٣	270.807	0	~	286.598	-1	2	G
10	276	365	0	٣	293.888	1	♥	269.943	9	4	1287,3962	8	4	10
11	27	491	9-	2	294.684		4.	269.076	7	7	288.192	9	4	11
12	27	15	-1	4	295.476		4	268.205	Н	٣				12
13	273	36	-2	4	296.264	ł	2	267.331	-2	7	289.771	0	7	13
14	27	52	-3	٣	297.048	1	89	266.454	0	~	290.555	0	٣	14
15	27	65	14	Ŋ	297.830		4	265.574	7	7	291.336	-5	2	15
16	27	91	7	4	298.608	1	9	264.691	~	7	292.115	0	7	16
17	27	82	-1	ည	299.383		e	263.805	٣	4	1292.8901	7	7	17
18	26	85	-4	മ	300,152	ł	7				293.661	ည	٣	18
19	26	84	-5	ഗ	300.921		2	262.021	0	7	294.429	0	7	19
20	26	81	-	4	301.683	,	9	261.125	-2	٣	295.194	7	7	20
21	26	73	-3	4	302.444		4	260.226	4	4	295.955	7	۳	21
22	26	63	-1	4	303.200	1	ß	259.323	4-	٣	296.713	0	7	22
23	26	48	9-	S	303.953		4	258.418	- 1	~	297.468	0	c	23
24	26	31	7	4	304.702	1	4	257.509	C	7	98.219	0	7	24
25	26	11	~	4	305.448		<u>س</u>	256.598	6	S	98.967	-2	٣	25
56	26	86	-3	7	306.190		က	255.682	7	~	99.715	58	24	56
27	26	59	9	S	306.928	J	٣	254.764	-1	7	00.454	4	3	27
28	26	28	4	0	307.663		٣	1253.8433	c	7	01.192	٣	~	28
59	25	94	4	m	308.394		4	252.918	-2	7	01.926	-1	7	59
30	25	99	7	٣	309.122		0	251.991	-	7	02.658	-1	7	30

31	32	33	34	35	36	37	38	39	40	4.1	42	43	4.1	45	46	47	48	49	20	51	52	53	54	55	99	57	58	59	9	61	62	63	64	65	99	<i>L</i> 9
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40.4311         -17         29         1154.7549         10         19         1140.4311         -17         29         1155.6088         25         19           1156.4604         8         10         1156.4604         7         10         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11         11 </td <td></td> <td></td> <td></td> <td>153.903</td> <td>9</td> <td></td> <td></td> <td></td> <td>153.903</td> <td>7</td> <td>7</td>				153.903	9				153.903	7	7
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32.9839         -3         9         1162.4650         15         17         1132.9839         -4         9         1162.4650         13         17         13         13         16         15         17         1132.1615         -12         12         1163.3260         3         11         1         1132.1615         -12         12         1163.3260         3         11         1         1163.3260         3         11         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1         1	33.808	-		61.602	-	9	133.8083 1	~	61,602	_	
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31.3421         1         8         1164.1881         -4         12         1131.3421         -1         8         1164.1881         -8         12         1         30.5232         2         10         1165.0526         -5         7         2         2         10         1165.0526         -5         7         2         2         10         1165.0526         -5         7         2         2         10         1165.0526         -5         7         2         2         10         1165.0526         -5         7         2         2         1165.0526         -5         7         2         2         1165.0526         -5         7         2         2         1166.0538         -6         11         1128.0730         1         7         1166.7838         -9         12         2         14         1128.0730         1         1167.6516         -5         14         2         2         1468.5198         6         9         1127.2589         2         9         1168.5198         -9         12         2         1126.4461         3         8         1169.3900         -15         8         1126.4461         3         8         1126.3300         12         112 <td< td=""><td>32.161</td><td>-1</td><td></td><td>163.326</td><td>S</td><td></td><td>132.1615 -1</td><td>7</td><td>163.326</td><td>1</td><td></td></td<>	32.161	-1		163.326	S		132.1615 -1	7	163.326	1	
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28.8879       1       1165.9193       17       14       1129.7032       -19       14       1165.9193       10       14       2         28.0730       1       1166.7838       0       12       1128.0879       -5       8       1166.7838       -9       12         28.0730       9       7       1167.6516       6       14       1128.0730       1       7       1167.6516       -5       14       2         27.2589       12       9       1168.5198       6       9       1127.2589       2       9       1168.5198       -9       9       2         26.4461       15       8       1126.4461       3       8       1169.3900       -3       8       2         25.6336       10       6       1170.2599       12       13       1126.4461       3       8       1169.3900       -3       8       2         24.8230       10       6       1177.02599       12       13       1124.8230       -9       8       11       1124.8230       -9       11       2       11       2       2       11       2       14       11       11       11       11       11       11	30.523		10	165.052	7		130.5232	–	165.052	2	
28.8879       1       8       1166.7838       -9       12         28.0730       9       7       1167.6516       6       14       1128.0730       1       7       1167.6516       -5       14       2         28.0730       9       7       1167.6516       6       14       1128.0730       1       7       1167.6516       -5       14       2         27.2589       12       9       1127.2589       2       9       1168.5198       -9       9       2         26.4461       15       8       1126.4461       3       8       1169.3900       -3       8       2         25.6336       10       6       1170.2599       12       13       1126.4461       3       8       1169.3900       -3       8       2         24.8230       10       6       1177.02599       12       13       1124.8230       -9       8       11       1124.8230       -9       8       11       1124.8230       -9       8       11       1172.0037       -22       11       2       11       2       11       2       11       2       11       2       11       2       11       2				165.919			129.7032 -1	~	165.919	0	
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27.2589       12       9       1168.5198       6       9       1127.2589       2       9       1168.5198       -9       9       2         26.4461       15       8       1169.3900       -3       8       2       8       1169.3900       -3       8       2         25.6336       10       6       1170.2599       12       13       1125.6336       -6       6       1170.2599       -12       13       2         24.8230       11       8       11       1122.0037       14       9       1172.0037       -22       9       1172.0037       -22       9       2       1172.0037       -22       9       2       1172.0037       -22       9       2       1172.0037       -22       9       2       11       2       2       11       1123.2045       -27       14       1172.0037       -22       9       2       11       2       2       11       2       2       11       2       2       11       2       11       2       11       2       2       11       2       11       2       2       11       2       11       2       11       2       11       11	28.073		7	167.651	9		128.0730		167.651	5 1	
26.4461       15       8       1126.4461       3       8       1169.3900       -3       8       2         25.6336       10       6       1170.2599       12       13       1125.6336       -6       6       1170.2599       -12       13       2         24.8230       11       8       11       113.1308       -2       11       2       11       2       11       2       11       2       2       11       2       2       11       2       3       11/2.0037       -2       9       11/2.0037       -2       9       2       11       2       3       2       11       2       3       2       11       2       3       3       2       11       2       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3       3<	27.258	7	σ	168.519	9	6	127.2589		168.519	6	
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